# EQUIPARTITIONING DESIGN METHOD OF A QUASI-PERIODICAL FOCUSING-ACCELERATING COUPLED SYSTEM

Fu Shinian, Guan Xialing\*, Fang Shouxian Institute of High Energy Physics, CAS, Beijing 100039, China \* China Institute of Atomic Energy, Beijing 102413, China

## Abstract

High-current proton linac sets a very severe requirement on the control of beam loss and emittance growth. Beam physics theory demonstrates that driven by the beam coherent instability; beam emittance grows due to the coupling of the strong space charge effect when an intense beam has different temperatures between different degrees of freedom. Therefore it is necessary to design an intense beam accelerator in line with the equipartitioning principle. However, this design becomes too tedious to reach full equipartitioning because a real proton linac is, in fact, a quasi-periodic coupling system. Taking the advantage of the code TRACE3-D[1], we modified it and added equipartitioning design function to it. It becomes a useful tool to readily realize equipartitioning design, with the assistance of PARMILA code[2]. This paper will introduce the modification to TRACE3-D code, and then illustrate its application to the design of a DTL linac section, which shows the importance of equipartitioning design in the intense beam linac.

# 1 INTRODUCTION

Beam loss is a major concern in designing an intense beam proton linac, because it may result in a radioactivity on the machine strong enough to exclude the manual access to the machine for the necessary maintenance during its long-term operation. It has been realized that the beam halo formation and emittance growth are important reasons for the beam loss. Beam physics study[3-4] reveals that the beam matching and equipartitioning design can effectively control the beam halo formation and emittance growth.

A proton linac uses different accelerating structure at different energy sections. It usually consists of RFQ, DTL, CCL and/or superconducting linac. Due to the increase of proton energy, the structures form quasiperiodic accelerating-focusing systems. PARMTEQ and PARMILA codes are commonly used for the designs of RFQ and the other accelerators, and TRACE3-D is frequently applied for the beam matching design. However, these codes were unable to achieve equipartitioning design. It is a tedious work to reach an equipartitioning design in a quasi-periodic system when one directly applies the equipartition theory without code assistance. RFQ design code PARMTEQ has been modified for the equipartitioning design and applied the

RFQ design of Frankfurt University. We modified TRACE3-D code in order for it to be capable of designing an equipartitioned linac. The modified code can readily achieve an equipartitioned linac design under the assistance of PARMILA code. In this paper, we will briefly introduce our modification to TRACE3-D code and use it for an example DTL linac to show its application. A major effort was put on keeping the beam matched while equipartitioned in the modification. Before describing the modification, the equipartition theory applied in the code will be outlined, as the basis of the code modification.

# 2 BASIS OF THE CODE MODIFICATOIN

Beam theory has demonstrated that beam should satisfy the equipartition condition in order to avoid the beam emittance growth and beam halo formation due to the coupling in different freedoms through strong space charge effect in an intense beam. For a matched beam in a uniform focusing channel, the condition is given in the following form:

$$\gamma_0 \frac{\mathcal{E}_{nx}}{\mathcal{E}_{nz}} \frac{z_m}{a} = 1, \qquad (1)$$

or, in another form, which we used in the code:

$$\frac{\varepsilon_{nx}}{\varepsilon_{nz}} \frac{k_x}{k_z} = 1. \tag{2}$$

In the formula,  $\epsilon_x$  and  $\epsilon_z$  are the beam normalized emittances in transversal and longitudinal directions, a and  $z_m$  the beam RMS size in the two directions,  $k_x$  and  $k_z$  the wave numbers with space charge.

Even though the formula is derived for uniform focusing channel, it is also applicable for a matched periodic focusing channel, if the wave numbers in equation (2) are recognized, under smooth approximation, as the ones of a periodic system. However, the real proton linac usually is not a real periodic system due to the acceleration. As the result, the wave numbers change along the linac. The variation of longitudinal wave number is usually first fixed by linac longitudinal dynamics. Then the focusing strength of the magnets is adjusted to achieve the equipartition condition. The zero-current wave number in longitudinal direction can be expressed in terms of the following form:

$$k_{z0} = \left[ -\frac{q}{mc^2} \frac{2\pi}{\lambda} \frac{1}{\beta_0^3 \gamma_0^3} \frac{1}{L} \sum_{i=1}^n E_{0i} L_i T_i \sin \phi_i \right]^{\frac{1}{2}}$$
 (3)

with q being the charge of the particle, m the mass of the particle, c the speed of the light in vacuum,  $\lambda$  the RF wavelength,  $\beta$  and  $\gamma$  the relativistic factors, L the total cell length in one period,  $E_{0i}$ ,  $L_i$ ,  $T_i$  and  $\phi_i$  the accelerating gradient, cell length, transit time factor and synchronous phase of the i-th cell in one period, n the total number of the cells in one period. The transversal phase advance of zero-current in one period  $\sigma_{x0}$  can be found by the trace of the transfer matrix M:

$$\cos \sigma_{x_0} = \frac{Tr(M)}{2} . \tag{4}$$
The wave number may

become a complex function of the parameters of a periodic system, and hence the equipartitioning design may turn out to be very tedious work. FODO channel of a DTL is a simple periodic system. The transversal transfer matrix of one period is:

$$M = D_4 Q_F D_3 G_2 D_2 Q_D D_1 G_1 \tag{5}$$

in which,

$$D_i = \begin{bmatrix} 1 & (L_i - l_Q)/2 \\ 0 & 1 \end{bmatrix}, G_i = \begin{bmatrix} 1 & 0 \\ k_i & 1 \end{bmatrix},$$

$$Q_{\scriptscriptstyle F} = egin{array}{ccc} \cos heta_{\scriptscriptstyle F} & rac{1}{\sqrt{\kappa_{\scriptscriptstyle F}}} \sin heta_{\scriptscriptstyle F} \ -\sqrt{\kappa_{\scriptscriptstyle F}} \sin heta_{\scriptscriptstyle F} & \cos heta_{\scriptscriptstyle F} \ \end{array} 
ight],$$

$$Q_{D} = \begin{bmatrix} \cosh \theta_{D} & \frac{1}{\sqrt{\kappa_{D}}} \sinh \theta_{D} \\ \sqrt{\kappa_{D}} \sinh \theta_{D} & \cosh \theta_{D} \end{bmatrix},$$

and 
$$\theta_{Q} = \sqrt{\kappa_{Q}} l_{Q}$$
,  $\kappa_{\varrho} = \frac{q}{mc} \frac{B_{\varrho}}{\beta_{\varrho} \gamma_{0}}$ ,  $B_{Q}$  is the magnetic

gradient of quadrupoles, l<sub>Q</sub> is the length of the quadrupole. Q stands for F or D, corresponding to focusing and defocusing magnets. The RF defocusing effect is represented by ki, with  $k_i = \frac{-\pi q E_{0i} L_i T_i \sin \phi_i}{mc^2 \lambda R^3 v^3}$ . To

realise equapartition design, one must select a properite  $B_0$  value to satisfy the contion of eq.(2).

## MODIFICATION OF TRACE3-D CODE

It is obvious from the last section that the equipartition design is a rather complicated procedure for a quasiperiodic system, even in a simple case like a FODO channel of a DTL linac. So a code is necessary to design an equipartitioned linac composed of hundreds of periods. We noticed that TRACE3-D code is suitable to design a matched linac and expanding its capability equipartitioning design is possible with the aid of PARMILA code. The longitudinal dynamic design of the PARMILA code outputs an input file for TRACE3-D, which gives out the quasi-periodic system and hence is a good start point. This output file can be further used for equipartitioning design by adjusting the magnetic gradient of focusing magnets in the modified TRACE3-D code. And more ever, TRACE3-D code itself includes the emittance growth due to the linear part of space charge effect and to RF field.

As the command K for periodic system in TRACE3-D was originally ineffective, it is adopted to enter the subroutine of equipartitioning design in the modified code. One can make choice of either equipartitioning design or a design of constant zero-current phase advance, which is also a kind of frequently used design method.

Because matching is the basis of the equipartitioning design, before equipartitioning design, one needs to design the input beam to be matched with the first period of the linac by making use of the matching design function of TRACE3-D. The modified code will revise the focusing field generated by PARMILA to satisfy the equipartition conditions, period by period. It is essential to maintain the beam always to be matched with focusing channel along the linac while changing the field of the magnets. So repeat may be necessary in some case. When a satisfied result is reached, the code will output a list of the focusing field for each magnet in the linac. It can then be put into PARMILA to replace the originally designed

# APPLICATION EXAMPLE

To show the application of the modified code in equipartition design, an example is presented in this section. Table 1 lists the major parameters of a section of a DTL linac used in the example.

Table 1 Major Parameters of the DTL Linac

Frequency f	324 MHz
Inject Energy W <sub>0</sub>	3 MeV
Output Energy W	13 MeV
Accelerating Gradient E <sub>0</sub>	2.5 MV/m
Synchronous Phase φ <sub>0</sub>	-30°
Number of Cells	51
Beam Current I	60 mA
Lattice	FODO
Nor. Emittance (RMS) $\epsilon_{nx}$	0.0931 mm-mrad
$\epsilon_{ny}$	0.0925 mm-mrad
$\epsilon_{nz}$	0.135 Deg-MeV

The equipartitioning factor  $\frac{\mathcal{E}_{nz}}{\mathcal{E}_{nx}}\frac{k_z}{k_x}$  is plotted in Figure

1 for the cases with a constant zero-current phase advance of 50 degree and with an equipartitioning condition. It is observed that the longitudinal temperature of the non-equipartitioned beam is three times higher than that of the transversal temperature at the entrance of the linac. For the equipartitioned beam the factor keeps unit almost all along the linac.

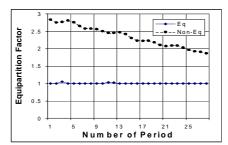


Figure 1: Equipartition factor along the linac for equipartitioned and non-equipartitioned beams (a constant zero-current phase advance of 50 degree in transversal direction).

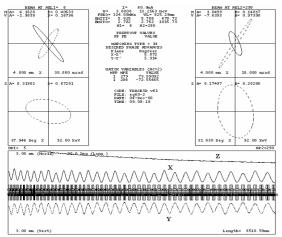


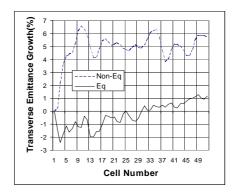
Figure 2: TRACE3-D output for the equipartitioned beam in the DTL. Up-left plot is the input beam phase spaces; up-right plot is the output beam phase spaces; bottom plot is the beam envelope in three directions.

TRACE3-D output the graph of the beam phase spaces and envelopes for the equipartitioned beam, as shown in Figure 2. The beam envelopes in the bottom plot exhibit a good matching of the beam with the focusing lattice.

To compare the emittance growth of the equipartitioned and non-equipartitioned beams, the magnetic field from TRACE3-D runs was copied into PARMILA input file. PARMILA runs with 10,000 macroparticles show the emittance variation in Figure 3. It is noticed that the transversal emittance of the equipartitioned beam increases only a little, several times less than that of the non-equipartitioned beam. On the other hand, in longitudinal direction, emittance growth for

equipartitioned beam is more than that of the nonequipartitioned beam. The results reflect the fact that the equipartitioning design prevents the temperature transfer from the longitudinal direction to the transversal direction.

The example in this section shows that the modified TRACE3-D code can be readily used for the design of equipartitioned beam in a quasi-periodic system with acceleration.



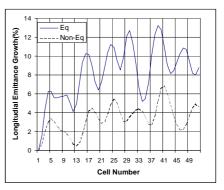


Figure 3: Beam emittance growth for equipartitioned and non-equipartitioned beam in transversal (up) and longitudinal (low) directions.

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